# Regional Climate Modeling over the Maritime Continent. Part II: New Parameterization for Autoconversion of Convective Rainfall

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#### ABSTRACT

This paper describes a new method for parameterizing the conversion of convective cloud liquid water to rainfall ("autoconversion") that can be used within large-scale climate models, and evaluates the new method using the Regional Climate Model, version 3 (RegCM3), coupled to the land surface scheme Integrated Biosphere Simulator (IBIS). The new method is derived from observed distributions of cloud water content and is constrained by observations of cloud droplet characteristics and climatological rainfall intensity. This new method explicitly accounts for subgrid variability with respect to cloud water density and is independent of model resolution, making it generally applicable for large-scale climate models. This work builds on the development of a new parameterization method for convective cloud fraction, which was described in Part I.

Simulations over the Maritime Continent using the Emanuel convection scheme show significant improvement in model performance, not only with respect to convective rainfall but also in shortwave radiation, net radiation, and turbulent surface fluxes of latent and sensible heat, without any additional modifications made to the simulation of those variables. Model improvements are demonstrated over a 19-yr validation period as well as a shorter 4-yr evaluation. Model performance with the Grell convection scheme is not similarly improved and reasons for this outcome are discussed. This work illustrates the importance of representing observed subgrid-scale variability in diurnally varying convective processes for simulations of the Maritime Continent region.

### 1. Introduction

The skill of large-scale climate models in reproducing the existing climate has improved significantly over many parts of the world, but simulations over the Maritime Continent region still contain substantial error (e.g., Yang and Slingo 2001; Neale and Slingo 2003; Dai and Trenberth 2004; Wang et al. 2007). Gianotti et al. (2012) showed that the Regional Climate Model, version 3.0 (RegCM3), coupled with the land surface scheme Integrated Biosphere Simulator (IBIS) also exhibited significant error over the Maritime Continent region with respect to rainfall, net radiation, latent and sensible heat fluxes, and evapotranspiration over the land surface. These errors are not unique but are consistent with other published climate modeling studies over the Maritime Continent region (e.g., Chow et al. 2006; Francisco et al. 2006; Martin et al. 2006; Wang et al. 2007). It was argued that the source of these errors resided primarily in the

atmospheric component (i.e., RegCM3) of the coupled model system, specifically with the simulation of convective cloud and rainfall processes (Gianotti et al. 2012).

In the Intergovernmental Panel on Climate Change Fourth Assessment Report (IPCC AR4), cloud feedbacks were identified as a primary reason for differences between models, with the shortwave impact of boundary layer and midlevel clouds making the largest contribution (Christensen et al. 2007). Of the wide range of time scales exhibited by clouds, variability on the diurnal scale is particularly important since it influences diurnal energetics of the atmosphere and surface (Bergman and Salby 1996). However, it was shown in Gianotti (2012) and Gianotti and Eltahir (2014, hereafter Part I) that the default parameterization of convective cloud fraction in RegCM3 did not simulate a reasonable diurnally varying convective-radiative feedback. This issue was addressed in Part I by introducing a new method for parameterizing convective cloud fraction.

However, the formation of convective clouds represents only one part of the diurnal cycle of convective cloud activity. The production of convective rainfall is a significant cloud dissipation mechanism in climate

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models and therefore is a critical component of convective cloud simulation. Most GCMs and RCMs use one or more associated equations to parameterize the suite of processes responsible for precipitation production, which is termed "autoconversion."

This work describes a new parameterization method for convective autoconversion and its implementation within RegCM3–IBIS. The new method explicitly accounts for subgrid variability in rainfall production and is independent of model resolution, making it generally applicable for use in large-scale climate models.

## 2. Parameterization methods for convective autoconversion

## a. Commonly used methods in large-scale climate models

Global and regional climate models (GCMs and RCMs) usually employ bulk microphysics schemes in which the hydrometeor representation is reduced to two variables—cloud droplets and precipitating particles—that are characterized by their mixing ratios (Geoffroy et al. 2008). Hence most autoconversion parameterization methods relate rainfall production to some dependence on cloud liquid water. Arguably the most commonly used of these methods are those developed by Kessler (1969), Arakawa and Schubert (1974), and Sundqvist et al. (1989).

In the scheme developed by Kessler (1969), the efficiency of autoconversion increases linearly with cloud water content but is zero below some threshold, that is,

$$\frac{dM}{dt} = k_1(m-a) \begin{cases} k_1 > 0 & \text{when } m > a \\ k_1 = 0 & \text{when } m < a \end{cases}, \quad (1)$$

where M is the rainwater mixing ratio, m is the cloud liquid water (CLW) mixing ratio,  $k_1 = 10^{-3} \, \text{s}^{-1}$ , and a is a CLW threshold value below which cloud conversion does not occur. Kessler (1969) noted that a reasonable threshold value would be  $1 \, \text{g m}^{-3}$ . This form of autoconversion was implemented in the Colorado State University GCM (Fowler et al. 1996) and in the fifthgeneration Pennsylvania State University–National Center for Atmospheric Research Mesoscale Model (MM5; Grell et al. 1994), as well as in the cloud-resolving model System for Atmospheric Modeling, version 6.3 (SAM 6.3; Khairoutdinov and Randall 2003).

Arakawa and Schubert (1974) used the following parameterization for the rainfall rate out of a subensemble of cumulus clouds at a given elevation *z*:

$$r(z) = C_0 l(z), \tag{2}$$

where  $C_0$  is a constant parameter and l(z) is the mixing ratio of CLW.

Lord (1982) proposed a similar autoconversion function that includes a dependence on the vertical flux of cloud water in the updraft, where  $C_0 = 0.002 \,\mathrm{m}^{-1}$ . This is the form of autoconversion is used within the Community Atmosphere Model, version 4.0 (CAM4.0; Neale et al. 2010).

In the model described by Sundqvist et al. (1989), the rate of release of both convective and stratiform precipitation is given by

$$P = c_0 m \left\{ 1 - \exp \left[ -\left(\frac{m}{b m_r}\right)^2 \right] \right\},\tag{3}$$

where  $1/c_0$  is the characteristic time scale for the conversion of cloud droplets into raindrops, b is cloud cover,  $m_r$  the within-cloud threshold value for cloud water, and m the gridcell value of cloud water (hence m/b is the within-cloud value of cloud water). The parameter  $m_r$  is assigned a value typical of individual cloud types, which is invariant to grid resolution. The values of  $c_0$  and  $m_r$  are modified by temperature to account for the more rapid autoconversion processes involving ice. This method of autoconversion was adopted by Tiedtke (1993) and is used in the National Aeronautics and Space Administration (NASA) Goddard Institute for Space Studies (GISS) GCM (Del Genio et al. 1996).

Some large-scale climate models also include prognostic droplet size and concentration variables. Those models may use autoconversion parameterizations based on the method of Manton and Cotton (1977), which depends on the mean concentration, radius, and collision frequency of cloud water droplets. Once the critical droplet radius  $r_{\rm crit}$  is prescribed and the cloud droplet number concentration  $N_d$  is known at a grid point, the threshold cloud liquid-water mixing ratio below which no autoconversion occurs can be estimated. While this method is more sophisticated than others, many models do not include droplet concentration and radius as prognostic variables.

The major drawback of all these autoconversion methods is the lack of subgrid variability in cloud water content and rainfall production. The methods rely on grid-mean values of the simulated variables and thresholds, and therefore rainfall is produced uniformly over a model grid cell.

GCM and RCM simulations can exhibit strong sensitivity to the parameters controlling autoconversion. For models that use a critical cloud water mixing ratio threshold, it is common to use a value considerably smaller than the within-cloud value of about 1 g kg<sup>-1</sup> suggested by Kessler (1969). For example, the Met Office

Unified Model uses a threshold of  $0.2\,\mathrm{g\,kg^{-1}}$ , the Colorado State University GCM uses  $0.25\,\mathrm{g\,kg^{-1}}$ , the NASA GISS GCM uses  $0.5\,\mathrm{g\,m^{-3}}$ , and the Scripps Institution of Oceanography GCM uses  $0.3\,\mathrm{g\,kg^{-1}}$  (Xu et al. 2005). The cloud-resolving model SAM 6.3 employs a threshold of  $1\,\mathrm{g\,kg^{-1}}$  (Blossey et al. 2007) because it runs at a scale sufficiently small (with a horizontal resolution on the order of  $1\,\mathrm{km}$ ) to resolve the natural variability in convective activity. For models that have a prognostic droplet radius and concentration, observations suggest that the critical drop size radius is  $10{-}15\,\mu\mathrm{m}$ , but GCMs typically use  $4.5{-}7.5\,\mu\mathrm{m}$  in order to get good simulated values of cloud liquid water (Rotstayn 2000; Geoffroy et al. 2008).

These lower, tuned parameter values could be interpreted as grid-mean thresholds and therefore comparable to the grid-mean simulated condensate. But because these thresholds are not functions of the actual subgrid variability in condensate, they must be manually tuned to account for different model resolutions, domains, and choice of parameterization schemes. Their use is therefore less than ideal.

Rotstayn (2000) added a new treatment of stratiform clouds to the Commonwealth Scientific and Industrial Research Organisation (CSIRO) GCM to simulate the subgrid-scale variability in stratiform autoconversion. The assumed subgrid moisture distribution from the model's condensation scheme was applied in each grid box to determine the fraction of the cloudy area in which the mean within-cloud mixing ratio exceeded a prescribed threshold, and autoconversion occurred only in that fraction of the grid cell (Rotstayn 2000). The total water mixing ratio was assumed to follow a symmetric triangular probability density function (PDF) about its grid-mean value. This method allowed the critical value of the threshold droplet radius to increase from 7.5  $\mu$ m in the old treatment to 9.3 µm in the new treatment, close to observations (Rotstayn 2000). This scheme improves the representation of subgrid-scale variability in autoconversion. But it requires a prognostic droplet concentration variable, which is not always available in large-scale climate models, and employs a somewhat arbitrary symmetric triangular PDF to describe the total water mixing ratio.

# b. Existing parameterization of convective autoconversion within RegCM3

The Grell convective parameterization scheme available within RegCM3 uses a method based on Arakawa and Schubert (1974) to calculate the convective rainfall produced at a given elevation, R(z):

$$R(z) = \frac{c_0 \Delta z}{1 + c_0 \Delta z} (q_u - q_{\rm rch}) M_u, \tag{4}$$

where  $c_0$  is 0.002,  $\Delta z$  is the depth of model layer,  $q_u$  is the water vapor mixing ratio of updraft,  $q_{\rm rch}$  is the water vapor mixing ratio of environment, and  $M_u$  is updraft mass flux. The quantity shown in parentheses is equivalent to the grid-mean CLW for that model layer.

Since there is no discussion of this formulation in Grell (1993) or the RegCM3 manual, it is assumed that the value of  $c_0$  is taken from Lord (1982). Because the vertical model layers are defined within RegCM3 to be thinner in the lower atmosphere and increase in thickness with increasing height above the surface, this autoconversion function results in a minimum efficiency of about 0.5 in the lower atmosphere and a maximum efficiency of about 0.8 at high altitudes.

The Emanuel scheme within RegCM3 uses a Kessler form of autoconversion to calculate the production of convective rainfall at a given elevation, R(z):

$$R(z) = \overline{\text{CLW}}(z) - \text{CLW}_T, \tag{5}$$

where  $\overline{\text{CLW}}(z)$  is grid-scale CLW within model layer z, and CLW<sub>T</sub> is the threshold value of CLW. By default, CLW<sub>T</sub> = 1.1 g kg<sup>-1</sup>.

The default threshold value is close to observed values of cloud density. However, it represents a point-scale threshold value, while the grid-mean cloud water mixing ratio  $\overline{\text{CLW}}$  is used to calculate simulated rainfall. Hence there is a mismatch of scales in this formulation within RegCM3. In the work presented in Gianotti (2012) and Part I, the threshold was set to  $0.25\,\mathrm{g\,kg^{-1}}$  to attempt a better matching of scales, but that kind of tuning should ideally be avoided.

### 3. New parameterization method for convective autoconversion

A new method has been developed to parameterize convective autoconversion. This method is in the spirit of the semiempirical approach described by Emanuel (1994), in which physical principles are used to constrain a system to a limited number of parameters, whose values are then related to observations.

As a first-order approximation, the long-term average efficiency of convective rainfall production is considered constant. Empirical relationships are then used to frame an expression for the time-mean autoconversion, as follows. When nonzero rainfall is simulated within a model grid cell, the fractional area of the grid cell that contains rainfall can be estimated using the relationship derived by Eltahir and Bras (1993):

$$\mu = \frac{\overline{R}}{R_{\text{olim}}},\tag{6}$$

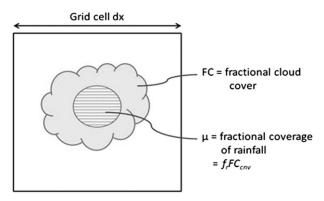


FIG. 1. Relationship between fractional cloud cover FC and fractional coverage of rainfall  $\mu$ .

where  $\mu$  is the fractional coverage of rainfall,  $\overline{R}$  is the grid-average simulated rainfall and  $R_{\rm clim}$  is the climatological rainfall intensity, which is the average rainfall intensity observed when there is nonnegligible rainfall. The fractional coverage of rainfall within a raining model grid cell may also be expressed as the portion of the cloudy area within a grid cell that produces rainfall:

$$\mu = f_r FC_{cnv}, \tag{7}$$

where  $f_r$  is the fraction of total cloudy area that is raining and FC<sub>cnv</sub> is the fractional coverage of convective cloud. This relationship is illustrated schematically in Fig. 1.

A new method for parameterizing convective cloud fraction is described by (Gianotti 2012; Part I)

$$FC_{cnv} = \frac{\overline{CLW}}{CLW_{clim}},$$
 (8)

where  $\overline{\text{CLW}}$  is the simulated grid-average value of CLW and CLW<sub>clim</sub> is the climatological cloud liquid water density (i.e., the mean observed cloud water content).

Combining the two expressions for the fractional coverage of rainfall yields

$$\frac{\overline{R}}{R_{\text{clim}}} = f_r \frac{\overline{\text{CLW}}}{\text{CLW}_{\text{clim}}}.$$
 (9)

The simulated grid-average rainfall can be represented as some fraction  $f_e$  of the simulated CLW, that is,

$$\overline{R} = f_{\rho} \overline{\text{CLW}},$$
 (10)

where  $f_e$  is the autoconversion efficiency, which can be expressed as

$$f_e = f_r \frac{R_{\text{clim}}}{\text{CLW}_{\text{clim}}}.$$
 (11)

The form of  $f_r$  is now required. In the physical world, if CLW is distributed according to a PDF, then the quantity  $f_r$  is equivalent to the portion of that PDF that exceeds some threshold for autoconversion, CLW<sub>T</sub>, that is,

$$f_r = \int_{\text{CLW}_T}^{\infty} f_{\text{CLW}} d\text{CLW}. \tag{12}$$

The observed distribution of CLW has been fitted to a lognormal distribution (Foster et al. 2006) and a Weibull distribution (Iassamen et al. 2009). Rainfall drop size distributions have been fitted to exponential (Marshall and Palmer 1948), lognormal (Rosenfeld and Ulbrich 2003), and gamma (Ulbrich 1983; Testud et al. 2001) distributions. For simplicity and in the absence of a welldefined PDF for CLW, an exponential distribution of CLW is used here (noting that other distributions may collapse to the exponential distribution with appropriate parameter choices). It should also be noted that a lognormal distribution for CLW was attempted in this work but no significant differences from an exponential distribution were found in the results. The exponential distribution is also advantageous since it can be integrated analytically a priori, increasing computational efficiency.

If the observed CLW is distributed exponentially with a mean of CLW<sub>clim</sub>, then

$$f_r = \int_{\text{CLW}_T}^{\infty} \frac{1}{\text{CLW}_{\text{clim}}} \exp\left(-\frac{\text{CLW}}{\text{CLW}_{\text{clim}}}\right) d\text{CLW}$$

$$= \exp\left(-\frac{\text{CLW}_T}{\text{CLW}_{\text{clim}}}\right), \tag{13}$$

which results in

$$f_e = \exp\left(-\frac{\text{CLW}_T}{\text{CLW}_{\text{clim}}}\right) \frac{R_{\text{clim}}}{\text{CLW}_{\text{clim}}}.$$
 (14)

This relationship represents the long-term mean climatological rainfall production efficiency, as shown in Fig. 2. However, at the point scale the autoconversion efficiency will be dynamic and spatially variable. It is desirable to capture this temporal and spatial variability within a model to improve the physical realism of the simulation. To achieve this, a point-scale expression for autoconversion can be used to derive an instantaneous autoconversion efficiency that accounts for the impact of subgrid-scale variability.

At the scale of a small parcel of convective cloud, the volume of rainfall will be a function of the cloud droplets that are large enough to fall, which can be represented as

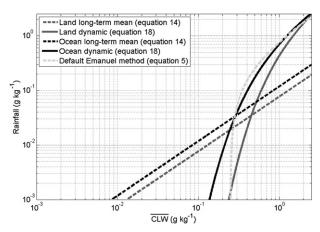


FIG. 2. Relationship between simulated grid-mean  $\overline{\text{CLW}}$  (g kg<sup>-1</sup>) and simulated grid-scale rainfall (g kg<sup>-1</sup>) calculated using the long-term mean autoconversion efficiency [Eq. (14)] compared to the instantaneous dynamic autoconversion efficiency [Eq. (18)], developed using typical values for land and ocean provided by observations (see Table 2). Note the log-log axes. Also shown for comparison is the default version of autoconversion within the Emanuel convection scheme in RegCM3 [Eq. (5)], using a threshold value of  $\text{CLW}_T = 0.25 \, \text{g kg}^{-1}$ .

$$R = \alpha (\text{CLW} - \text{CLW}_T)^n \text{ for } \text{CLW} > \text{CLW}_T,$$
 (15)

where R is point-scale rainfall, CLW is the amount of CLW within the cloud parcel, CLW<sub>T</sub> is the point-scale threshold of CLW that must be exceeded to produce rainfall, and n is the parameter that represents the degree of linearity of the autoconversion process. The coefficient  $\alpha$  (it is assumed that  $\alpha = 1$ ) ensures that the units remain dimensionally consistent for cases when  $n \neq 1$ .

This point-scale relationship follows the Kessler form and would duplicate the Kessler method if n = 1. However, the generalized form in which  $n \neq 1$  accounts for the many complexities involved in converting cloud droplets to rainfall, such as the effects of turbulent mixing and variability in cloud condensation nuclei size and abundance. Such nonlinear processes have been proposed as explanations for the difference between theory and observations of rainfall production (Emanuel 1994; Khain et al. 2000; Stephens and Haynes 2007).

If the convective CLW within a model grid cell is distributed with the PDF  $f_{\text{CLW}}$ , then  $\overline{R}$  can be written as

$$\overline{R} = \alpha \int_{\text{CLW}_{T}}^{\infty} (\text{CLW} - \text{CLW}_{T})^{n} f_{\text{CLW}} d\text{CLW}. \quad (16)$$

Within a model grid cell, if the simulated distribution of convective CLW  $f_{\text{CLW}}$  is also assumed to follow an exponential distribution with a mean value of  $\overline{\text{CLW}}$  (the simulated instantaneous grid-mean value of CLW), then

$$\overline{R} = \alpha \overline{\text{CLW}}^n \Gamma(n+1) \exp\left(-\frac{\text{CLW}_T}{\overline{\text{CLW}}}\right),$$
 (17)

where  $\Gamma$  is the gamma function, which can be evaluated via lookup tables.

This expression can now be used to derive a dynamic, grid-scale measure of the autoconversion efficiency:

$$f_e = \alpha \overline{\text{CLW}}^{n-1} \Gamma(n+1) \exp\left(-\frac{\text{CLW}_T}{\overline{\text{CLW}}}\right).$$
 (18)

This function requires specification of two parameter values:  $CLW_T$  and n, assuming  $\alpha=1$ . It is desirable to have these values constrained by observations. Unfortunately, observations of CLW are limited to non-precipitating retrievals and hence there are no direct observations of autoconversion. But observations have been made of the radius required for a cloud droplet to convert into a precipitating raindrop and of the droplet number concentration at the time that the critical threshold radius is breached. An approximate value of  $CLW_T$  can be estimated from these observations, as follows:

$$CLW_T = \frac{4}{3}\pi\rho N_{d,crit}r_{crit}^3,$$
 (19)

where  $\operatorname{CLW}_T$  is threshold cloud water content (g m<sup>-3</sup>),  $\rho$  is liquid water density (g m<sup>-3</sup>),  $N_{d,\text{crit}}$  is droplet concentration corresponding to critical radius (m<sup>-3</sup>), and  $r_{\text{crit}}$  is the critical radius for rainfall formation of cloud droplets (m).

Table 1 lists observations of the critical droplet radius and droplet concentration in continental and maritime clouds. Since the islands in the Maritime Continent often experience biomass burning, it seems reasonable to adopt  $r_{\rm crit} \approx 8\,\mu{\rm m}$  and  $N_d \approx 700\,{\rm cm}^{-3}$ . Using these values in Eq. (19) leads to  ${\rm CLW}_{T,{\rm land}} \approx 1.5\,{\rm g\,m}^{-3}$ . Over ocean, it seems reasonable to adopt  $r_{\rm crit} \approx 14\,\mu{\rm m}$  and  $N_d \approx 60\,{\rm cm}^{-3}$ . This leads to  ${\rm CLW}_{T,{\rm ocean}} \approx 0.7\,{\rm g\,m}^{-3}$ . But it also seems sensible to set  ${\rm CLW}_T > {\rm CLW}_{\rm clim}$ , since the average observed cloud density is likely to be nonprecipitating. Given that the observed value of  ${\rm CLW}_{\rm clim,ocean} \approx 0.7\,{\rm g\,m}^{-3}$  (Table 1), it seems reasonable to set  ${\rm CLW}_{T,{\rm ocean}} \approx 0.75\,{\rm g\,m}^{-3}$ .

It now remains to evaluate the parameter n, which could be attempted in multiple ways. One method is to constrain the instantaneous autoconversion efficiency by the long-term mean efficiency. This requires values for  $R_{\rm clim}$  and  ${\rm CLW_{clim}}$ .

Data from Singapore's Changi airport meteorological station (described in Gianotti et al. 2012), the Tropical Rainfall Measuring Mission (TRMM)'s 3B42 product (described in section 4), and Table 1 of Eltahir and Bras

Cumulus cloud type	Quantity	Value	Reference Rosenfeld and Lensky (1998)		
Continental, general	Liquid water content	$0.1-3 \mathrm{g}\mathrm{m}^{-3}$			
_	Liquid water content	$1  \mathrm{g  m}^{-3}$	Rogers and Yau (1989)		
	Critical droplet radius, $r_{\rm crit}$	9–10 μm	Brenguier et al. (2000)		
	Droplet concentration, $N_d$	Median = $228 \mathrm{cm}^{-3}$ , third quartile = $310 \mathrm{cm}^{-3}$	Squires (1958)		
Continental,	Critical droplet radius, $r_{crit}$	$3-8 \mu m$	Reid et al. (1999)		
biomass burning	Critical droplet concentration, $N_d$	$3000\mathrm{cm}^{-3}$	Reid et al. (1999)		
Maritime	Liquid water content	$0.25-1.3\mathrm{gm^{-3}}$	Rangno and Hobbs (2005)		
	Liquid water content	$0.4-1.2\mathrm{gm}^{-3}$	Warner (1955)		
	Critical droplet radius, $r_{crit}$	15 μm	Kubar et al. (2009)		
	Critical droplet radius, $r_{\rm crit}$	$12-15 \mu \text{m}$	Rangno and Hobbs (2005)		
	Critical droplet radius, $r_{\rm crit}$	$13-14  \mu \text{m}$	Brenguier et al. (2000)		
	Critical droplet concentration, $N_d$	$60  \text{cm}^{-3}$	Kubar et al. (2009)		
	Critical droplet concentration, $N_d$	$70{\rm cm}^{-3}$	Rangno and Hobbs (2005)		

 $Median = 45 cm^{-3}$ 

TABLE 1. Observations used to constrain new autoconversion function.

(1993) suggest  $R_{\rm clim,land} \approx 4.85 \, {\rm mm \, h^{-1}}$  and  $R_{\rm clim,ocean} \approx 3.5 \, {\rm mm \, h^{-1}}$  for the Maritime Continent.

Droplet concentration,  $N_d$ 

To provide a direct comparison between CLW and rainfall, the climatological rainfall rates  $R_{\rm clim}$  (mm h<sup>-1</sup>) can be converted to mixing ratios (i.e., the mass of rainfall falling through a unit volume of air), by assuming an exponential size distribution for raindrops as in Marshall and Palmer (1948) and by using the functions provided by Rosenfeld and Ulbrich (2003). These yield  $R_{\rm clim,land} = 0.32 \, {\rm g \, m^{-3} \, h^{-1}}$  and  $R_{\rm clim,ocean} = 0.24 \, {\rm g \, m^{-3} \, h^{-1}}$ .

The data presented in Table 1 suggest that  $CLW_{clim,land} \approx 1.2 \, \mathrm{g \, m^{-3}}$  and  $CLW_{clim,ocean} \approx 0.7 \, \mathrm{g \, m^{-3}}$ . Using these values in Eq. (14) leads to a mean autoconversion efficiency of  $f_e \approx 0.1$  over both land and ocean.

Values of the simulated grid-scale  $\overline{\text{CLW}}$  are now required to evaluate the dynamic autoconversion efficiency using Eq. (18). With both the Grell and Emanuel convection schemes in RegCM3, typical instantaneous values of  $\overline{\text{CLW}}$  were approximately 0.5 g m<sup>-3</sup> over land and 0.3 g m<sup>-3</sup> over ocean. Model output is not shown for brevity. Using these values in Eq. (18) leads to n = 0.9 over land and n = 0.94 over ocean.

Another possible way to constrain parameter n is to set the efficiency equal to 1 at some upper threshold value for  $\overline{\text{CLW}}$ . It is noted that the same parameter values of n=0.9 over land and n=0.94 over ocean can be derived by assuming that  $f_e=1$  when  $\overline{\text{CLW}}=2.5\,\text{g m}^{-3}$  over land and  $2\,\text{g m}^{-3}$  over ocean, close to the highest observed values of CLW (see Table 1).

Hence the following expressions are obtained for the grid-mean convective rainfall:

$$\overline{R} = \overline{\text{CLW}}^{0.9} 0.9618 \exp\left(-\frac{0.0015}{\overline{\text{CLW}}}\right) \text{ for land}$$
 (20a)

and

$$\overline{R} = \overline{\text{CLW}}^{0.94} 0.9761 \exp\left(-\frac{0.00075}{\overline{\text{CLW}}}\right) \text{ for ocean,}$$
(20b)

**Squires** (1958)

where  $\overline{\text{CLW}}$  is the simulated grid-mean CLW (kg kg<sup>-1</sup>). These relationships are illustrated in Fig. 2. The figure shows that rainfall production using the dynamic method is much more efficient than the long-term mean autoconversion efficiency for high  $\overline{\text{CLW}}$  values but less efficient at low  $\overline{\text{CLW}}$  values.

RegCM3 uses only a single prognostic variable, liquid water mixing ratio, but in reality the autoconversion efficiency is strongly dependent upon the type of hydrometeor. In particular, the conversion of cloud ice crystals to precipitation is known to be much faster than the conversion of cloud liquid droplets (Rogers and Yau 1989). To account for this difference, the Emanuel convection scheme scales  $CLW_T$  in temperatures below a certain threshold, as follows:

$$CLW_{T,actual} = CLW_T$$
 for  $T \ge 0$ °C, (21a)

 $CLW_{T,actual}$ 

$$= \text{CLW}_T \left( 1 - \frac{T}{-55^{\circ}\text{C}} \right) \quad \text{for} \quad -55^{\circ}\text{C} \le T \le 0^{\circ}\text{C} ,$$
(21b)

and

$$CLW_{T,actual} = 0$$
 for  $T \le -55^{\circ}C$ . (21c)

This increases the autoconversion efficiency in cold clouds where ice crystals would be expected to dominate. This same scaling is included in the new autoconversion method presented here, but a detailed treatment of

Parameter	Land	Ocean	Notes
$R_{ m clim}$	$4.85 \mathrm{mm}\mathrm{h}^{-1}, \\ 0.32 \mathrm{g}\mathrm{m}^{-3}\mathrm{h}^{-1}$	$3.5 \mathrm{mm}\mathrm{h}^{-1}, \\ 0.24 \mathrm{g}\mathrm{m}^{-3}\mathrm{h}^{-1}$	Data from Changi airport, Singapore (Gianotti et al. 2012), TRMM 3B42 product (see section 3c) (Table 1 of Eltahir and Bras 1993). Unit conversion used functions from Marshall and Palmer (1948) and Rosenfeld and Ulbrich (2003).
$\text{CLW}_{\text{clim}}$	$1.2\mathrm{gm}^{-3}$	$0.7\mathrm{gm}^{-3}$	Estimated from observations in Rangno and Hobbs (2005), Rogers and Yau (1989), Rosenfeld and Lensky (1998), and Warner (1955).
$CLW_T$	$1.5\mathrm{gm}^{-3}$	$0.75\mathrm{gm}^{-3}$	Derived from observations in Brenguier et al. (2000), Kubar et al. (2009), Rangno and Hobbs (2005), Reid et al. (1999), and Squires (1958)
n	0.94	0.9	Derived by matching instantaneous autoconversion efficiency to long-term mean efficiency.
CLW	$2.5 \mathrm{g}\mathrm{m}^{-3}$	$2\mathrm{gm^{-3}}$	Taken from model output (not shown).

TABLE 2. Values of parameters used in derivation of new convective autoconversion parameterization.

ice-phase precipitation should be a priority for future work.

Table 2 summarizes the parameter values used in this derivation. However, it is noted that these values are chosen from within ranges of observed values (Table 1). Therefore it would be reasonable for the model user to choose parameter values within these ranges to explore the impact of different parameter choices on model performance.

It is noted that if a value of n=1 is assumed with the new formulation, and the corresponding value of  $\operatorname{CLW}_T$  is derived using observed values of  $R_{\text{clim}}$  and  $\operatorname{CLW}_{\text{clim}}$  for the Maritime Continent region, a value of  $\operatorname{CLW}_T \simeq 0.3 \, \mathrm{g \, m}^{-3}$  is obtained. This is similar to the "tuned" values of  $\operatorname{CLW}_T$  used in some GCMs including that used in Part I. Indeed, Fig. 2 shows the similarity between the dynamic autoconversion and the tuned version of the default Emanuel method. The generalized form presented here, where n=1 is not assumed a priori, removes the need to tune a grid-mean value.

The new formulation for autoconversion presented here has significant advantages over the default forms that exist within RegCM3: 1) it explicitly recognizes subgrid variability in CLW and how that variability affects the grid-scale conversion process; 2) only two parameters have to be specified, one dependent upon the other, which are constrained by observations; and 3) it can be applied consistently across convection schemes, bringing added realism to a model independently of user choices.

# **4. Simulations using the new convective autoconversion parameterization**

#### a. Model description

This work uses RegCM3–IBIS as described in Gianotti et al. (2012), including the subgrid explicit moisture SUBEX scheme (Pal et al. 2000) for resolvable,

large-scale clouds and precipitation, the choice of Grell (Grell 1993) with Fritsch-Chappell (F-C; Fritsch and Chappell 1980) or Arakawa-Schubert (Grell et al. 1994) closures, and Emanuel (Emanuel 1991; Emanuel and Živković-Rothman 1999) convective parameterization schemes. Further details of the developments and description of RegCM3 are available in Pal et al. (2007).

Giorgi et al. (2012) describes upgrades that were incorporated into the more recent version RegCM, version 4 (RegCM4), which was made publicly available in 2011. RegCM4 does not contain any upgrades to convective parameterization relevant to this work.

Some modifications were made to the boundary layer parameterization scheme (Holtslag et al. 1990; Holtslag and Boville 1993) within RegCM3, the simulation of large-scale clouds within the PBL, soil thermal conductivity, and ocean surface roughness (see Gianotti 2012 for details). The vertical limit on simulated cloud cover was also extended to permit clouds up to an altitude of about 16 km. The combination of these modifications resulted in small reductions to the land and ocean surface latent heat fluxes, a substantial reduction in the PBL height, removing an overestimation bias in the PBL height over land at night, and removal of egregious nighttime low-level large-scale clouds over land (Gianotti 2012). However, none of these modifications significantly impacted the simulation of radiative fluxes or rainfall.

#### b. Experimental design

Two sets of simulations were run to evaluate the performance of the new autoconversion parameterization within RegCM3–IBIS. In all simulations described below, the new parameterization was implemented as follows.

Both the Grell and Emanuel convection schemes in RegCM3 simulate rainfall within a model layer using a function of the form shown in Eq. (10), that is, as the product of some autoconversion efficiency and the grid-mean CLW within that layer. In the Grell scheme, the efficiency is a function of the grid-mean updraft mass flux  $M_u$  and two constant parameters (model layer thickness and a coefficient). Note that the quantity shown in parentheses in Eq. (4) is equivalent to the grid-mean CLW. In the Emanuel scheme, the efficiency is a function of the CLW threshold, CLW<sub>T</sub>. Hence in both convection schemes, the new autoconversion efficiency function shown in Eq. (18) simply replaces the existing efficiency function such that the rainfall calculation remains of the form shown in Eq. (10). Since both convection schemes calculate the amount of condensate (i.e., CLW) resulting from the grid-mean water vapor mixing ratio and saturation vapor pressure, the new autoconversion parameterization can be implemented consistently across convection schemes.

The first simulations were shorter and were used to evaluate the impact of the new autoconversion method relative to other model modifications made by these authors. Included in this set are simulations in which the planetary boundary layer region has been modified as described above [labeled Mod; see Gianotti (2012) for details], simulations incorporating both boundary layer modifications and the new parameterization of convective cloud fraction shown in Eq. (8) (labeled New; see Part I for details), and simulations combining the new convective autoconversion with all other modifications (labeled Auto). Simulations began on 1 July 1997 and ended 31 December 2001. The first 6 months of output were ignored to allow for spinup. The remaining 4 years of simulation (1998–2001) were used for evaluation of model performance.

A second set of two simulations was run to validate the performance of the newly modified RegCM3–IBIS over a longer time period. One simulation uses the default version of RegCM3–IBIS [labeled Def; for details see Gianotti et al. (2012)]; the other simulation incorporates all the model modifications mentioned here (i.e., as in Auto). These simulations used the Emanuel

TABLE 3. List of parameters used in SUBEX and their default values (based on Pal et al. 2000).

Parameter	Land	Ocean
Cloud formation threshold RH <sub>min</sub>	0.8	0.9
Maximum saturation RH <sub>max</sub>	1.01	1.01
Autoconversion rate $C_{ppt}$ (s <sup>-1</sup> )	$5 \times 10^{-4}$	$5 \times 10^{-4}$
Autoconversion scale factor $C_{acs}$	0.65	0.3
Accretion rate $C_{acc}$ (m <sup>3</sup> kg <sup>-1</sup> s <sup>-1</sup> )	6	6
Raindrop evaporation rate $C_{\text{evap}} \left[ (\text{kg m}^{-2} \text{s}^{-1})^{-1/2} \text{s}^{-1} \right]$	$2\times10^{-5}$	$2 \times 10^{-5}$

convection scheme, beginning on 1 July 1982 and ending 31 December 2001. The first 6 months of output were ignored to allow for spinup. The remaining 19 years of simulation (1983–2001) were used for model evaluation. These years were chosen so that the same datasets could be used for lateral boundary conditions as for the 4-yr simulations, with similar observational datasets, for consistency.

Table 3 presents the parameter values used in the SUBEX routine in this study. The default values were based on the original version of the scheme, as described in Pal et al. (2000). It should be noted that although the large-scale SUBEX routine and the convective parameterization schemes operate independently within the RegCM3 structure, each routine can significantly impact the other by altering grid-scale variables such as temperature and water vapor. Therefore the model user should be cognizant of how changes introduced to the large-scale scheme affect convection, including the impact of the new parameterization presented here, and vice versa.

For all simulations, the model domain (Fig. 3) was centered along the equator at 115°E, used a normal Mercator projection, and spanned 95 grid points meridionally and 200 grid points zonally, with a horizontal resolution of 30 km. The simulations used 18 vertical sigma levels, from the ground surface up to the 50-mb level. The land surface scheme was run every 120 s, twice the model time step.

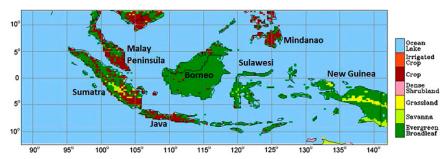


FIG. 3. Model domain showing vegetation classification used for IBIS. The domain has been sized such that it covers the generally accepted extent of the Maritime Continent region. Major islands are labeled.

TABLE 4. Varying characteristics of simulations used in study. The names are used to reference each simulation in the text.

Simulation name	Convection scheme	Modified PBL region*	Convective cloud fraction	Convective autoconversion	Evaluation period
GFC-Mod	Grell with F-C	Yes	Default	Default	1998-2001
GFC-New	Grell with F-C	Yes	New**	Default	1998-2001
GFC-Auto	Grell with F-C	Yes	New**	New	1998-2001
EMAN-Mod	Emanuel	Yes	Default	Default	1998-2001
EMAN-New	Emanuel	Yes	New**	Default	1998-2001
EMAN-Auto	Emanuel	Yes	New**	New	1998-2001
EMAN-Def_Long	Emanuel	No	Default	Default	1983-2001
EMAN-Auto_Long	Emanuel	Yes	New**	New	1983-2001

<sup>\*</sup> Refers to simulation of PBL height, ocean surface roughness, soil heat flux, and large-scale clouds within the PBL region, as described in Gianotti (2012).

Sea surface temperatures (SSTs) were prescribed using the National Ocean and Atmospheric Administration (NOAA) optimally interpolated SST (OISST) dataset, which is available at 1° × 1° resolution and at a weekly time scale beginning in 1982 (Reynolds et al. 2002). Topographic information was taken from the United States Geological Survey's Global 30-arc-s elevation dataset (GTOPO30), aggregated to 10 arc min (United States Geological Survey 1996). Vegetation biomes were based on the potential global vegetation dataset of Ramankutty (1999), modified to include two extra biomes for inland water and ocean as described in Winter et al. (2009). In all simulations presented, RegCM3–IBIS was run only with static vegetation.

Soil properties, including albedo and porosity, were determined based on the relative proportions of clay and sand in each grid cell. Sand and clay percentages were taken from the Global Soil Dataset, which has a spatial resolution of 5 min (Global Soil Data Task, IGDP-DIS 2000). Soil moisture, temperature, and ice content were initialized using the output from a global  $0.5^{\circ} \times 0.5^{\circ}$  resolution 20-yr offline simulation of IBIS as described in Winter et al. (2009).

The 40-yr European Centre for Medium-Range Weather Forecasts (ECMWF) Re-Analysis (ERA-40) dataset, available from September 1957 to August 2002 (Uppala et al. 2005), was used to force the boundaries in all simulations. The exponential relaxation technique of Davies and Turner (1977) was used with both datasets.

Eight simulations in total are presented in this study. Table 4 summarizes the different characteristics of these simulations and lists the names used to reference each simulation throughout the text.

#### c. Comparison datasets

Solar radiation is compared to the NASA/Global Energy and Water Cycle Experiment (GEWEX) Surface Radiation Budget (SRB) dataset release 3.0 (Stackhouse et al. 2011), made available by the NASA Langley Research Center Atmospheric Sciences Data Center. SRB is available at 3-hourly intervals from 1983 onward on a  $1^{\circ} \times 1^{\circ}$  global grid. Data were interpolated to the model domain for direct comparison.

Observations of latent heat flux over land are taken from field studies of evapotranspiration (ET) in the Maritime Continent, as described in Table 4 of Gianotti et al. (2012). For all simulations, it is assumed that the fluxes measured by these field studies apply over the entire evaluation period. Over ocean, latent and sensible heat flux observations are from the Woods Hole Oceanographic Institution (WHOI) global dataset of ocean evaporation (Yu et al. 2008), which is available from 1958 onward on a  $1^{\circ} \times 1^{\circ}$  global grid.

Rainfall from the 4-yr simulations is compared to the TRMM Multisatellite Precipitation Analysis (TMPA)  $0.25^{\circ} \times 0.25^{\circ}$  resolution 3B42 product (described in Huffman et al. 2007), available from 1998 onward and referenced in this work simply as TRMM. Relative proportions of convective and large-scale rain are taken from Mori et al. (2004), who used the 2A25, 2A12, and 2B31 TRMM products to describe the climatological convective versus stratiform rainfall split over Indonesia.

Total rainfall from the 19-yr simulations is compared to the Global Precipitation Climatology Project (GPCP) version 2.2 combined precipitation dataset, available at monthly resolution from 1979 to the present on a 2.5° × 2.5° global grid (Adler et al. 2003). GPCP data are provided by the NOAA Office of Oceanic and Atmospheric Research/Earth System Research Laboratory Physical Sciences Division (NOAA/OAR/ESRL PSD) from their website (http://www.esrl.noaa.gov/psd/). It is assumed that the convective and large-scale rainfall fractions obtained from the TRMM data also apply over the 19-yr evaluation period; these fractions are used to separate the total rainfall from GPCP into convective and large-scale components.

<sup>\*\*</sup> As described in Gianotti (2012) and Part I.

Fractional cloud cover from the 4-yr simulations is compared to the International Satellite Cloud Climatology Project (ISCCP) stage D2 product (Rossow et al. 1996), made available by NASA. This provides the fractional cloud cover at three elevations: high (50-440 mb), middle (440–680 mb), and low (680–1000 mb), at 280-km resolution and monthly time scales. Only low and midlevel clouds are shown here, since these are the regions most affected by convective clouds and rainfall. ISCCP data were averaged over the period 1998–2001 for comparison to the model. Model output was aggregated to the same horizontal grid as the ISCCP data. To match up the vertical resolution, the model output was aggregated in the vertical assuming random overlap of clouds between layers, using layers 2–8 (roughly 760– 1000 mb) for the low clouds and layers 9–12 (roughly 450-700 mb) for the middle clouds.

#### 5. Results and discussion

#### a. Cloud fraction

Figures 4 and 5 show the time-mean low and midlevel cloud cover. The figures compare all six of the 4-yr (1998–2001) simulations presented in this study to the ISCCP data and illustrate the first-order effect of the new autoconversion parameterization. The figures show that GFC-Auto and EMAN-Auto simulations (see Table 4 for simulation names and characteristics) have significantly less low and middle cloud cover than GFC-New and EMAN-New. The difference is especially dramatic for EMAN-Auto. For both convection schemes, the previous overestimation bias in low cloud cover relative to ISCCP is significantly reduced in the Auto simulations, and the middle cloud is reduced to levels comparable to that of the Mod simulations.

These results suggest that 1) the new autoconversion method is significantly more efficient at producing rainfall and dissipating cloud cover than the default methods in RegCM3, and 2) the reasonable match between the default convective cloud cover (in GFC-Mod and EMAN-Mod) and ISCCP was the result of two compensating deficiencies: in creating the clouds (addressed in GFC-New and EMAN-New) and in dissipating the clouds via rainfall (addressed in GFC-Auto and EMAN-Auto). Hence the default cloud cover was reasonably close to ISCCP resulting from what could be considered a good combination of tuning. With the new formulations for convective cloud cover and autoconversion, the same cloud cover is produced but with much more physical realism. This has important ramifications for other aspects of the simulation, described further below.

Figures 6 and 7 present the mean diurnal cycle of the vertical structure of cloud cover averaged over land grid cells within the domain, for all six of the 4-yr (1998–2001) simulations presented in this study. The ocean cloud fraction is not shown since it does not exhibit a significant diurnal cycle. The figures show that the midlevel cloud cover that was prominent in GFC-New and EMAN-New is almost entirely removed in GFC-Auto and EMAN-Auto, again demonstrating the increased efficiency of the new autoconversion method in regions of high CLW. It is noted that the temperature falls below freezing at approximately 5-km elevation over this domain (not shown), and hence the midlevel cloud is mostly warm (i.e., liquid phase) and not in the presence of significant quantities of ice.

Figures 4 and 5 indicate that the time-mean cloud cover compared to ISCCP is not very different between the Mod and Auto simulations. But Figs. 6 and 7 clearly show that the diurnal cycles are shifted between these simulations such that the maximum low cloud cover occurs in the afternoon in GFC-Auto and EMAN-Auto rather than at night, as was the case for GFC-Mod and EMAN-Mod. This significantly impacts the simulation of radiative fluxes. Cloud fraction results from the 19-yr simulations are not shown here for brevity but were very similar to the 4-yr simulation results.

### b. Radiative and turbulent heat fluxes

Figures 8 and 9 show the average diurnal cycle of incoming solar radiation (insolation) reaching the surface over land and ocean, comparing the SRB observations to all six of the 4-yr (1998–2001) simulations. Mean daily values are given in parentheses.

With GFC-Auto, insolation was slightly reduced relative to GFC-Mod and GFC-New, but not enough to remove the overestimation bias compared to SRB. Since the simulated low and midlevel fractional cloud cover in GFC-Auto compared well to ISCCP, it is likely that the overestimation of insolation is the result of insufficiently dense cloud cover, particularly at high elevations (not shown). This outcome is the result of weak convective updraft mass flux simulated by the Grell scheme, which is discussed in detail further below.

Over both land and ocean, EMAN-Auto simulates more insolation than EMAN-New but less than EMAN-Mod, resulting in a diurnal cycle of insolation that matches well to SRB. It is especially encouraging that the peak insolation is well simulated, both in magnitude and timing. These results illustrate the impact of reproducing a reasonable diurnal cycle of cloud cover on radiative fluxes.

The results from the 4-yr simulations using the Emanuel convection scheme indicate that combining the new parameterizations for convective cloud fraction and

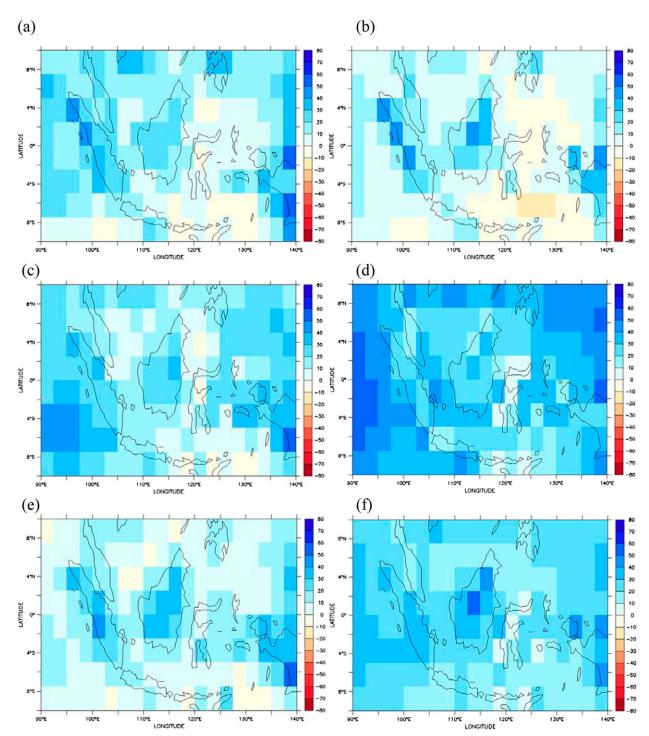


FIG. 4. Average low cloud fraction for 1998–2001: simulation minus ISCCP for (a) GFC-Mod, (b) EMAN-Mod, (c) GFC-New, (d) EMAN-New, (e) GFC-Auto, and (f) EMAN-Auto. Color bar indicates fractional coverage of grid cell.

autoconversion result in good simulation of both cloud cover and insolation. To validate that these model improvements hold over the 19-yr evaluation period and that other performance metrics are similarly improved, Table 5 summarizes the average daily surface radiative

and turbulent heat fluxes for the period 1983–2001 for the EMAN-Def\_Long and EMAN-Auto\_Long simulations compared to the observations.

Over land, EMAN-Def\_Long significantly overestimates insolation at the earth's surface. This occurs

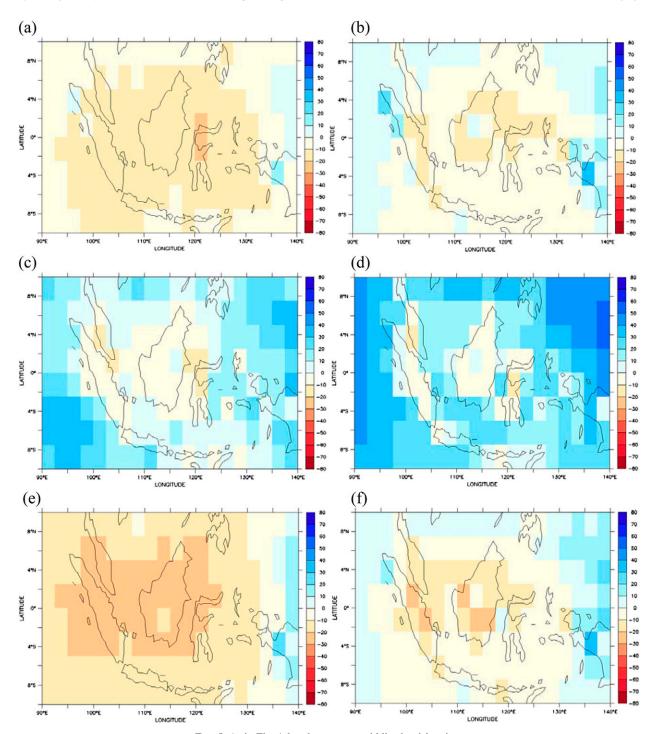


FIG. 5. As in Fig. 4, but for average middle cloud fraction.

despite good simulation of the planetary albedo and propagates into error in the simulated net radiation and the latent and sensible heat flux (LH and SH). But with EMAN-Auto\_Long, simulated insolation matches very well to SRB. There is a small underestimation of planetary

albedo. EMAN-Auto\_Long has a small underestimation of net radiation, primarily because of a combination of small errors in the simulated longwave radiation. Both EMAN-Def\_Long and EMAN-Auto\_Long show underestimation error in SH, although the bias is reduced

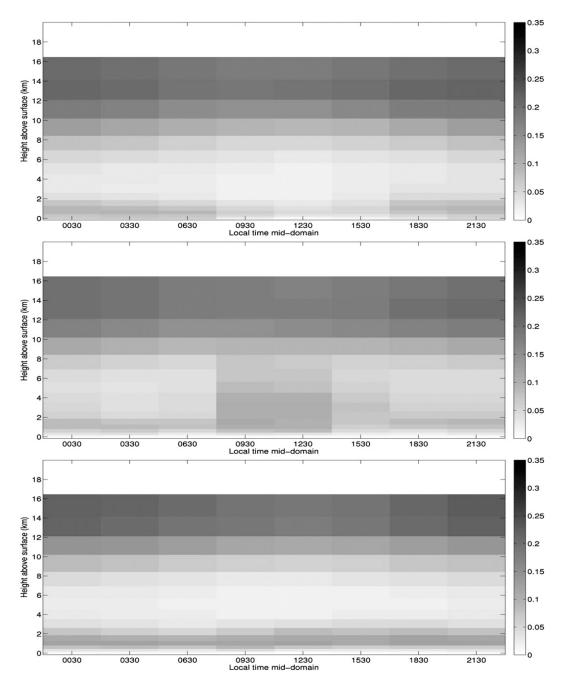


FIG. 6. Diurnal cycle of cloud cover averaged over land grid cells within the model domain for period 1998–2001 using simulations (top) GFC-Mod, (middle) GFC-New, and (bottom) GFC-Auto. The *x*-axis labels indicate the time of the middle of each 3-h output window, with respect to local time in the center of the model domain. To represent the *y* axis on a linear scale, the vertical extent of each model layer was assigned a single value of cloud cover, as provided by the model output. In reality, a smoother profile with less abrupt vertical variability would be expected.

in EMAN-Auto\_Long. This error most likely stems from the IBIS land surface scheme, not the atmospheric component of RegCM3.

Over ocean, the results from EMAN-Auto\_Long also show considerable improvement in the simulation of

insolation, with a very good match to observations and removal of the significant overestimation bias that was present in EMAN-Def\_Long. There is a small underestimation of net radiation with EMAN-Auto\_Long, again primarily because of small errors in the longwave

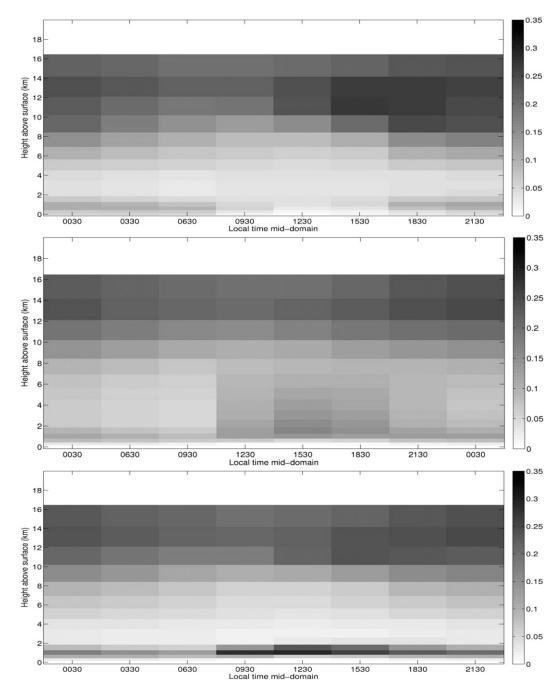


FIG. 7. As for Fig. 6, but using simulations (top) EMAN-Mod, (middle) EMAN-New, and (bottom) EMAN-Auto.

radiation. The simulated overestimation of LH and underestimation of SH in EMAN-Def\_Long is not addressed in EMAN-Auto\_Long. It is considered that this is because SSTs are fixed in these simulations, and hence turbulent heat fluxes over ocean exhibit little sensitivity to the surface net radiation. However, if an ocean model were coupled to the new RegCM3–IBIS model, it is expected that simulated turbulent fluxes over

the ocean would show improvement relative to the default version of RegCM3.

Therefore these results indicate that improvements to the simulated diurnal cycle of cloud cover and insolation propagate into improvements in the mean radiative and turbulent heat fluxes over longer time periods, when using the Emanuel convection scheme.

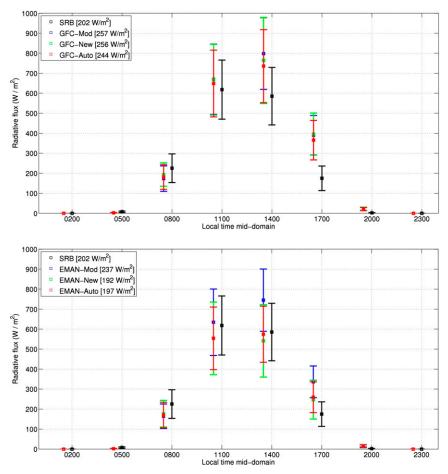


FIG. 8. Diurnal cycle of incoming solar radiation averaged over land for period 1998–2001 for SRB observations and simulations using (top) Grell with Fritsch–Chappell and (bottom) Emanuel convection schemes. Square symbol indicates the mean value; error bars indicate  $\pm 1$  standard deviation.

#### c. Rainfall

Table 6 shows the mean total, convective, and large-scale rainfall volumes for simulations evaluated over the period 1998–2001. The results show divergent behavior between the two convection schemes in response to the new autoconversion function.

GFC-Auto shows a reduction in convective rainfall compared to GFC-Mod and GFC-New, leading to an underestimation of convective and total rainfall over both land and ocean that was not present in the default version of RegCM3, resulting in worse model performance than the original version of this scheme (see Gianotti et al. 2012). This result seems inconsistent with the decrease in cloud cover and insolation that resulted from the new autoconversion method and is somewhat counterintuitive. The reasons for this are discussed in the next section.

EMAN-Auto shows an increase in convective rainfall relative to EMAN-New, but produces less rainfall than

in EMAN-Mod. This result is consistent with the shift in diurnal timing of cloud cover resulting from the new convective cloud fraction and associated change to insolation. The simulated convective and total rainfall volumes produced by EMAN-Auto are a reasonable match to TRMM over both land and ocean, showing considerable improvement on the previous simulations using the Emanuel scheme. These results indicate that the new autoconversion function, combined with the new convective cloud fraction, leads to improved physical realism throughout the simulation when using this scheme.

Table 6 also displays the total rainfall volumes and estimated convective and large-scale rainfall fractions produced by the 19-yr simulations. The results confirm that the new methods lead to improved model performance with respect to rainfall using the Emanuel scheme.

### d. Implications of using different convection schemes

Differences in model performance between the Grell and Emanuel convection schemes result from very

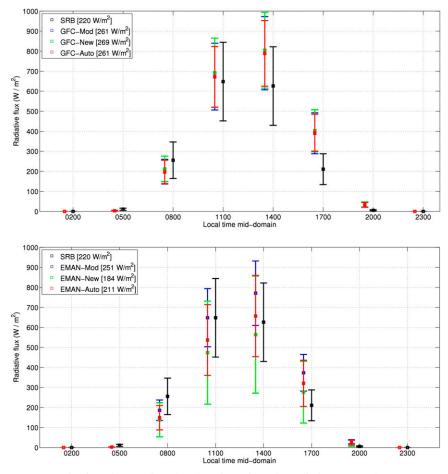


FIG. 9. As in Fig. 8, but for diurnal cycle of incoming solar radiation averaged over ocean.

different approaches to simulating convective updraft mass flux. The work presented here indicates that it is important to consider the implications of each scheme in the context of the Maritime Continent and other regions where convection is a dominant feature.

In both schemes, the volume of convective rainfall received at the surface is the product of the updraft mass flux and the amount of precipitable water produced in each model layer through the vertical column, minus any

reevaporation of rainfall that occurs between the cloud and the ground. Model testing showed that the amount of CLW produced from the available moisture within each grid cell was similar between the two schemes, as were the rates of reevaporation. Hence the primary cause of the difference in convective rainfall between the Grell and Emanuel schemes was the simulation of updraft mass flux.

In the Emanuel scheme, the updraft mass flux is a function of the degree of instability at the lifting

TABLE 5. Average daily surface radiation and turbulent heat fluxes over the period 1983–2001 for (top) land and (bottom) ocean cells within the domain. All radiative and turbulent fluxes are in units of W m<sup>-2</sup>.

Product/simulation	$SW_{dn}$	$SW_{up}$	$SW_{net}$	Land albedo	Planetary albedo	$LW_{dn}$	$LW_{up}$	$R_N$	LH	SH
					Land					
Observations	207	28	180	13%	47%	410	452	138	95	43
EMAN-Def_Long	226	32	195	14%	48%	412	460	147	126	22
EMAN-Auto_Long	208	30	179	14%	44%	406	457	128	100	29
					Ocean					
Observations	225	14	211	6%	44%	419	466	164	102	8
EMAN-Def_Long	271	17	254	6%	43%	412	471	194	118	5
EMAN-Auto_Long	226	15	211	7%	46%	419	471	159	119	5

TABLE 6. Average daily rainfall over land and ocean for observations and simulations evaluated over (top) 1998–2001 and (bottom) 1983–2001. All values are in units of mm day <sup>-1</sup>.

Product/simulation		Land average		Ocean average			
	Total	Convective	Large-scale	Total	Convective	Large-scale	
			1998	-2001			
TRMM	8.6	5.4 (63%)	3.2 (37%)	7.0	4.0 (57%)	3.0 (43%)	
GFC-Mod	10.9	4.4 (40%)	6.5 (60%)	8.5	4.1 (48%)	4.4 (52%)	
GFC-New	9.2	4.1 (45%)	5.1 (55%)	7.3	5.5 (75%)	1.8 (25%)	
GFC-Auto	5.9	3.1 (52%)	2.8 (48%)	4.8	2.9 (61%)	1.9 (39%)	
EMAN-Mod	16.8	9.9 (59%)	6.9 (41%)	6.7	3.8 (57%)	2.9 (43%)	
EMAN-New	10.3	5.0 (49%)	5.3 (51%)	7.9	3.5 (44%)	4.4 (56%)	
EMAN-Auto	9.9	5.4 (55%)	4.5 (45%)	6.1	3.7 (61%)	2.4 (39%)	
		, ,	1983	-2001	, ,	· · ·	
Observations*	7.2	4.5 (63%)	2.7 (37%)	5.7	3.2 (57%)	2.5 (43%)	
EMAN-Def_Long	10.3	6.6 (64%)	3.7 (36%)	3.9	3.2 (82%)	0.7 (18%)	
EMAN-Auto_Long	7.3	4.5 (62%)	2.8 (38%)	4.0	2.6 (65%)	1.4 (35%)	

<sup>\*</sup> Total rainfall volumes are taken from GPCP data; convective and large-scale rainfall fractions are taken from TRMM data as described in Mori et al. (2004) and applied to the GPCP rainfall volume.

condensation level (LCL), as measured by the difference in density temperature between a lifted air parcel and the environment at that time step (Emanuel 1991). This makes convection with the Emanuel scheme very sensitive to changes in the near-surface environment, such that convective rainfall directly reflects lower atmospheric instability.

Updraft mass flux in the Grell scheme is based on the cloud work function, which is an integral measure of the buoyancy energy available for convection and provides the necessary closure for the scheme. The updraft mass flux is derived based on the assumption that the cloud work function is in quasi-equilibrium, such that the time rate of change of the total cloud work, which is a function of both the large-scale variables and the modification of the environment because of the cumulus cloud ensemble, is approximately zero (Arakawa and Schubert 1974). For both the large-scale and cumulus cloud components, the cloud work is a function of the moist static energy in the updraft compared to the downdraft. The result of this assumption is that updraft mass flux and convective rainfall increase (decrease) if the updraft strength increases (decreases) relative to the downdraft.

This study illustrates the impact of these different approaches to updraft mass flux. When rainfall efficiency is increased using the Emanuel scheme, cloud cover decreases and insolation increases, creating more instability in the lower atmosphere and consequently stronger updraft mass flux. This creates a positive feedback loop that leads to increased rainfall. However, when rainfall efficiency is increased using the Grell scheme, this dries out the atmosphere and strengthens the downdraft, leading to a decrease in the net convective mass flux and producing less total rainfall at the surface.

In developing their cumulus parameterization, Arakawa and Schubert (1974) used a time-integrated updraft mass flux that represented the mean cloud mass flux over the life time of a cumulus cloud ensemble, taken to be the time scale for moist convective adjustment. This time scale was shown to be on the order of  $10^3$ – $10^4$ s (from 30 min to 3 h), significantly smaller than the time scale of the large-scale processes that create the instability driving convection, which is on the order of  $10^5$ s (about 1 day) (Arakawa and Schubert 1974).

The convective adjustment time scale was intended to be virtually instantaneous in the sense that it is basically the same as the computational time step for implementing the physics (Arakawa 2004). It is noted that the GCMs into which the original Arakawa and Schubert (1974) scheme was originally implemented were of resolution 100–200 km, with a computational time step on the order of hours, suitable for assuming an almost instantaneous adjustment due to convection.

However, an RCM running at the scale of tens of kilometers typically has a time step of a few minutes. Hence a single computational time step only represents a fraction of the cloud lifetime, substantially less than that required to assume "instantaneous" convective adjustment. Therefore the time-averaged mass flux used in the Arakawa and Schubert (1974) scheme is not considered appropriate for simulations that require a description of the time evolution of convection. The Emanuel scheme, on the other hand, achieves adjustment toward quasi-equilibrium during the time integration of explicitly formulated transient processes (Arakawa 2004). Therefore it is suggested that the Grell scheme is not suitable for RCM simulations but the Emanuel scheme is well suited to RCM simulations,

especially over regions of strong convection and given appropriate radiative forcing as shown in this work.

#### 6. Summary

This paper describes a new parameterization for convective autoconversion that can be used within large-scale climate models, and evaluates the new method using the coupled regional model RegCM3–IBIS. The new method is derived from observed distributions of cloud water content and is constrained by observations of cloud droplet characteristics and climatological rainfall intensity. This method explicitly accounts for subgrid variability with respect to cloud cover and cloud water density. It is spatially and temporally variable and independent of model resolution.

Simulations over the Maritime Continent using the Emanuel convection scheme showed improvement in model performance across a broad range of metrics, including cloud cover, radiative fluxes, and rainfall. The wet bias simulated by the default Emanuel scheme over land was greatly reduced without creating a dry bias over the ocean. The land to ocean rainfall ratio simulated by EMAN-Mod was 2.5:1, while the ratio simulated by EMAN-Auto was 1.6:1—a significant improvement compared to the observed ratio of 1.2:1 shown by TRMM.

The improvements in model performance with the Emanuel scheme held not only over a 4-yr evaluation period but also over a 19-yr validation period. If the model modifications were the result of parameter tuning over the initial 4-yr period, then longer simulations would not necessarily exhibit reasonable performance. Additionally, all the model modifications tested here were aimed at improving the representation of diurnal-scale processes, but model performance was improved over much longer time scales. These results support the assertion that poor representation of the diurnal cycle is a major source of error in climate simulations over the Maritime Continent.

Simulations with the Grell convection scheme produced counterintuitive results of increased rainfall efficiency but decreased total rainfall production. The results illustrate some of the limitations of the quasi-equilibrium theory for moist convective adjustment, as embodied in the Grell scheme, when used in a model with a relatively small time step over a region where convection is a dominant feature.

There are likely to be other important feedbacks of these new formulations for cloud cover and auto-conversion that are not documented here, such as sensitivity to different large-scale forcings. When coupled to an ocean model, instead of being forced with SSTs as in this study, it is also likely that the new formulations

for cloud cover and autoconversion will produce very different results than the default version of the model. Ongoing work will evaluate the model performance over regions other than the Maritime Continent and coupled to an ocean model.

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